



Polymer surface properties control the function of heavy meromyosin in dynamic nanodevices



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ARTICLE INFO

Article history:

Received 1 June 2016

Received in revised form

16 August 2016

Accepted 18 August 2016

Available online 20 August 2016

Keywords:

Molecular motors

Protein adsorption

Nanodevices

ABSTRACT

The actin-myosin system, responsible for muscle contraction, is also the force-generating element in dynamic nanodevices operating with surface-immobilized motor proteins. These devices require materials that are amenable to micro- and nano-fabrication, but also preserve the bioactivity of molecular motors. The complexity of the protein-surface systems is greatly amplified by those of the polymer-fluid interface; and of the structure and function of molecular motors, making the study of these interactions critical to the success of molecular motor-based nanodevices. We measured the density of the adsorbed motor protein (heavy meromyosin, HMM) using quartz crystal microbalance; and motor bioactivity with ATPase assay, on a set of model surfaces, i.e., nitrocellulose, polystyrene, poly(methyl methacrylate), and poly(butyl methacrylate), poly(tert-butyl methacrylate). A higher hydrophobicity of the adsorbing material translates in a higher total number of HMM molecules per unit area, but also in a lower uptake of water, and a lower ratio of active per total HMM molecules per unit area. We also measured the motility characteristics of actin filaments on the model surfaces, i.e., velocity, smoothness and deflection of movement, determined via in vitro motility assays. The filament velocities were found to be controlled by the relative number of active HMM per total motors, rather than their absolute surface density. The study allowed the formulation of the general engineering principles for the selection of polymeric materials for the manufacturing of dynamic nanodevices using protein molecular motors.

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1. Introduction

Mechanical work in biological nanosystems is performed by a variety of force-generating protein motors, such as myosins, kinesins and dyneins (Spudich, 2011; Vale, 2003; Veigel and Schmidt, 2011), the former being responsible for muscle contraction (A.F. Huxley and Niedergerke, 1954; H. Huxley and Hanson, 1954). In the “gliding geometry” motility assay developed in the late 1980s, whole myosin molecules (Kron and Spudich, 1986; Uyeda et al., 1991), or the part of the mechano-enzyme containing its working arms, e.g., heavy meromyosin (HMM) (Uyeda et al., 1991), or even the end of an arm, i.e., the S1 unit (Toyoshima et al., 1987; Uyeda et al., 1991), are adsorbed on a surface. Provided that the upper solution contains sufficient adenosine triphosphate

(ATP), the fluorescently-labelled actin filaments will be propelled by the surface-bound motors, sliding randomly on the surface, thus allowing the easy observation and quantification of motility characteristics using simple optical fluorescence microscopy set-ups and imaging software.

While the “gliding geometry” motility assay has been used extensively in fundamental studies of molecular motor function (Holzbaur and Goldman, 2010), from an applications perspective, their planar architecture and the ability of motor proteins to transport nano-scale cargo at speeds that are orders of magnitude higher than those associated with molecular diffusion (Nitta and Hess, 2005), are very attractive features for dynamic nanodevices (Bakewell and Nicolau, 2007; Fulga et al., 2009; Kinbara and Aida, 2005). Consequently, proof-of-concept motor-powered nanodevices have been proposed for biosensing (Agarwal et al., 2009; Martinez-Neira et al., 2005; Van Zalinge et al., 2012), biondiagnostics (Fischer et al., 2009; Korten et al., 2010), transport at nano- (Bull et al., 2005), and micro-scale (Limberis and Stewart, 2000), microfluidic pumping (Bull et al., 2005) and recently

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biocomputation (Nicolau et al., 2016).

Because the manufacturing of these devices must use materials that are both suitable for micro/nanofabrication, and also preserve the motor bioactivity, various materials have been assessed, e.g., methacrylate polymers (Nicolau et al., 1999; Riveline et al., 1998; Suzuki et al., 1997), polyurethane (Clemmens et al., 2003a), plasma polymerised poly(ethylene oxide) (Clemmens et al., 2003b), polyelectrolytes (Jaber et al., 2003), commercial photoresists (Bunk et al., 2003a, 2003b; Clemmens et al., 2004; Hiratsuka et al., 2001; Moorjani et al., 2003), and silane-functionalized surfaces (Bunk et al., 2005; Sundberg et al., 2003). Despite this rather large amount of empirical information, as well as several studies focused on the fundamentals of the motor protein-surface interactions (Albet-Torres et al., 2007a; Katira et al., 2007, 2009; Van Zalinge et al., 2012), there are still many uncertainties regarding the impact of surfaces on motor function, in particular regarding polymers, which are the preferred materials for inexpensive devices, due to the coupled complexities of the polymer and the protein systems.

To this end, to progress on the selection of materials for future dynamic nanodevices using actin-myosin system, we studied the relationship between the physico-chemical properties polymeric surfaces, in particular their hydrophobicity and polymer network structure, on one side; and the surface density of molecular motors and the preservation of their motility, on the other.

2. Materials and methods

2.1. Polymer surface coating

Superclean nitrocellulose (NC) was purchased from Ernest F. Fullam, Inc. (Latham, NY). Polystyrene (PS), poly(methyl methacrylate) (PMMA), poly(butyl methacrylate) (PBMA), poly(*tert*-butyl methacrylate) (PtBMA) and hexamethyldisilazane (HMDS) were purchased from Aldrich Chemicals. The selection of polymeric surfaces aims to reach a reasonably large range of properties related to motility assays and related devices: (i) nitrocellulose is the standard polymer for motility assays, but unfit for the fabrication of devices due to its flammability; (ii) polystyrene is the material

of choice for the plastic utensils in molecular biology and biochemistry, but is rarely used for motility assays; (iii) PMMA is the material of choice for polymer-based microfluidics devices, and has been also used with good results for motility assays; (iv) it would be useful to compare PMMA with other more hydrophobic acrylates, but with very different properties related to water uptake, i.e., high, and low T_g, for PtBuMA and PBMA, respectively. Finally, while silane polymers, e.g., poly(di methyl siloxane), PDMS, are used indeed for microfluidics devices, the release of mono-/oligo-mers is extremely toxic for motility assays.

The chemical structures for all model surfaces are presented in Fig. 1.

Glass coverslips were cleaned by sonication in 70% ethanol, dried in a stream of N₂, primed with HMDS, spin-coated with one of the polymer solutions, i.e., NC (1% w-v in amyl acetate), PS (2.5% w-v in propylene glycol monomethyl ether acetate, PGMEA), PMMA (2% w-v PGMEA), PBMA (1% w-v toluene), or PtBMA (2% w-v in PGMEA) at 3600 rpm, then soft-baked at 85 °C for three hours. The concentrations of solvents in the polymer solutions have been optimised to obtain a viscosity that leads to a smooth film during the spin coating.

The hydrophobicities of the polymer-coated surfaces were determined by contact-angle measurements using deionized water ($R > 18.2 \text{ M}\Omega$) and Krüss contact-anglemeter (DSA10Mk2). The reported values are averages of ten different readings for each surface.

The measurement of the viscoelastic properties of the polymers in thin films used an advanced commercial quartz microbalance (QCM) system (QCM-Z500, from KSV Instruments). This system allowed the measurement of the impedance spectrum, thus providing both the frequency and the bandwidth, addressing up to 11th harmonics. The measurements have been performed sequentially, in a step-wise manner, i.e., first on the bare dry polymer surfaces, then on surfaces interfaced with the buffer solution. A complete description of the QCM equipment, measurement protocols and associated theoretical background is presented in the [Supplementary information section](#).

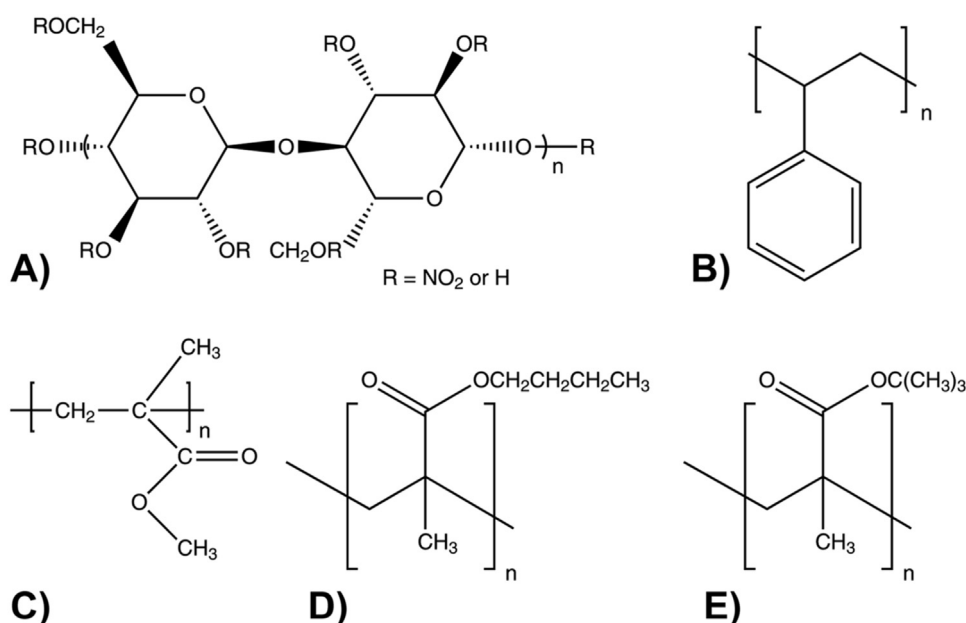


Fig. 1. Chemical structures of surfaces and polymers tested for in vitro actomyosin motility support. A: nitrocellulose (NC); B: poly(styrene) (PS); C: poly(methyl methacrylate) (PMMA); D: poly(butyl methacrylate) (PBMA); E: poly(*tert*-butyl methacrylate) (PtBMA).

2.2. Protein preparation

Myosin and G-actin were extracted and purified from rabbit skeletal back muscle using previously reported methods (Barden and Dosremedios, 1984; Carsten and Mommaerts, 1963; Margosian and Lowey, 1982; Spudich and Watt, 1971). HMM was prepared by proteolytic cleavage of myosin (20 mg mL⁻¹ in 0.5 M KCl, 50 mM KPO₄ pH 6.5) with 0.07 mg mL⁻¹ α -chymotrypsin. Prior to use in motility studies, G-actin was dissolved in 1 mM DTT to a concentration of 2 mg mL⁻¹ and clarified by centrifugation (40,000 rpm for 60 min). F-actin was simultaneously polymerised and fluorescently labelled by incubation of equal molar quantities of G-actin and AlexaFluor 488 phalloidin (Molecular Probes) in 4 mM imidazole pH 7, 2 mM MgCl₂, 0.1 mM EGTA, 3 mM NaN₃, 1 mM DTT.

2.3. Motility assays

Actomyosin motility assay procedures were modified from Sellers et al. (Sellers et al., 1993). The motility flow cells were constructed by sealing two parallel edges of a polymer-coated coverslip to a standard glass microscope slide using double-sided tape (approximately 100 μ m thickness) as a spacer. HMM was diluted to 0.1 mg mL⁻¹ in 10 mM MOPS pH 7.2, 20 mM KCl, 5 mM MgCl₂, 0.1 mM EGTA, 10 mM DTT (Buffer A, which is classically used for myosin-based motility assays, e.g., Sellers et al., 1993), introduced into the flow cell and allowed to attach on the surface for 2 min. The cell was then flushed sequentially with the following (all in Buffer A): (i) 1 mg mL⁻¹ BSA, (ii) 5 μ M unlabelled F-actin, (iii) 1.5 mM MgATP, and (iv) 5 nM labelled actin. The motility was initiated by the addition of 1 mM MgATP in Buffer A. A photo-bleaching inhibition solution, consisting of 2.5 mg mL⁻¹ glucose, 0.1 mg mL⁻¹ glucose oxidase (Sigma G-7016) and 0.02 mg mL⁻¹ catalase (Sigma C-100), was also added.

2.4. Motility data analysis

The motility was observed at room temperature (23–25 °C) on an inverted microscope (Olympus IX71) using a PlanApo 100x oil objective (Olympus), epifluorescence optics (FITC filter set), and mercury light source. Images were recorded every 150 ms for 60 s using a Coolview FDI high-resolution camera (Photronics Science Ltd.) controlled by Image-Pro Plus software (Ver. 5.0, Media Cybernetics). The manual object tracking features of Image-Pro Plus were used to determine the coordinates of actin filament heads in consecutive frames, from which the motility statistics were derived.

Replicate trials for each surface were conducted, with all filaments within one quadrant (representing a 44 \times 33 μ m² area) of the field of view being evaluated, resulting in 20–32 filaments being analysed for each surface.

The step velocity was calculated as the distance travelled by the filament head divided by the time between sampled frames. The step velocity variation was calculated as the difference between two consecutive velocities. The deflection angle was calculated for each combination of three consecutive points as the difference between (i) the direction of filament head movement between two consecutive points; and (ii) the direction of filament head movement between the next two consecutive points. The absolute value of the directional difference was calculated, such that the angle of the directional change was always less than 180°.

2.5. Measurement of the surface density of adsorbed HMM

The same advanced QCM system (QCM-Z500, KSV Instruments) was used to measure the amount of adsorbed protein mass, and to

assess the dynamics of viscoelastic properties of the thin films comprising proteins on, or embedded in polymers. The measurements used similar protocols as for the measurement of viscoelastic properties of the dry, and buffer-soaked polymer films. A complete description of the QCM equipment, measurement protocols and associated theoretical background is presented in the [Supplementary information section](#).

2.6. Measurement of HMM activity

The amount of active HMM bound to polymer-coated coverslips was estimated by comparing NH₄-EDTA ATPase activities in flow cells with those of known amounts of HMM in solution (all at 25 °C). In solution, ATPase activities were determined colorimetrically based on the rate of inorganic phosphate release (Taussky and Shorr, 1953) by known concentrations of HMM in ATPase buffer (25 mM imidazole hydrochloride, 0.2 M NH₄Cl, 9 mM EDTA, pH 7.4). The reaction was started by addition of 5 mM ATP and stopped by the addition of an equal volume of Taussky-Shorr reagent (1% [w: v] ammonium molybdate, 5% [w: v] ferrous sulphate heptahydrate in 0.5 M H₂SO₄). ATPase activities of surface immobilized HMM were determined in flow cells of HMM under standard motility assay conditions (see further). After equilibration of HMM in Buffer A, the flow cell was rinsed with ATPase buffer and the reaction started by the introduction of 5 mM ATP in ATPase buffer. After set reaction times, chosen at 5, and 10 min, respectively, the flow cell contents were removed by perfusion with 4 volumes (~200 μ L) ATPase buffer (no ATP), the total volume was brought to 1 mL, and the reaction stopped by the addition of an equal volume of Taussky-Shorr reagent. In order to ensure a full and stable consumption of ATP, we worked with HMM concentrations that are double than those reported elsewhere (Guo and Guilford, 2004) as limiting the rate of ATP consumption.

3. Results and discussion

3.1. Complex relationship between the surface properties and motility characteristics

The characteristics of the motility of cytoskeletal filaments on surface immobilized motor proteins in *in vitro* motility assays is a function of (i) the characteristics intrinsic of the motor system, e.g., ATP cycling time, power stroke length, strong binding time (Fulgata et al., 2009; Veigel and Schmidt, 2011), and (ii) the assay-related characteristics, e.g., the density of motor proteins on the surface, their activity, and ATP concentration (Albet-Torres et al., 2007a; Nicolau et al., 2007a). Aside of the above elements, the characteristics of the adsorbing surface in controlling the surface density of motors and their surface-controlled bioactivity, but the measurement of *both* these parameters has been very rarely reported. For instance, early experimental and theoretical studies (Harris and Warshaw, 1993; Rivelino et al., 1998; Uyeda et al., 1990) examined the correlation between the velocity of actin filaments as a function of the surface density of *active* motors (and their intrinsic parameters), as estimated from ATPase assays. Moreover, the analysis used a single type of surface, most frequently nitrocellulose, but overlooking the surface density of denatured, inactive motor molecules. Alternatively, the impact of the surface properties on the motor bioactivity was indirectly estimated by comparing the average actin filaments velocity on motor protein immobilized on various surfaces (Bunk et al., 2003a; Jaber et al., 2003; Nicolau et al., 2007b, 1999; Sundberg et al., 2003). Finally, the amount of surface-immobilized motors was measured using QCM (Albet-Torres et al., 2010, 2007b; Van Zalinge et al., 2012), but the concomitant report of the still active motors,

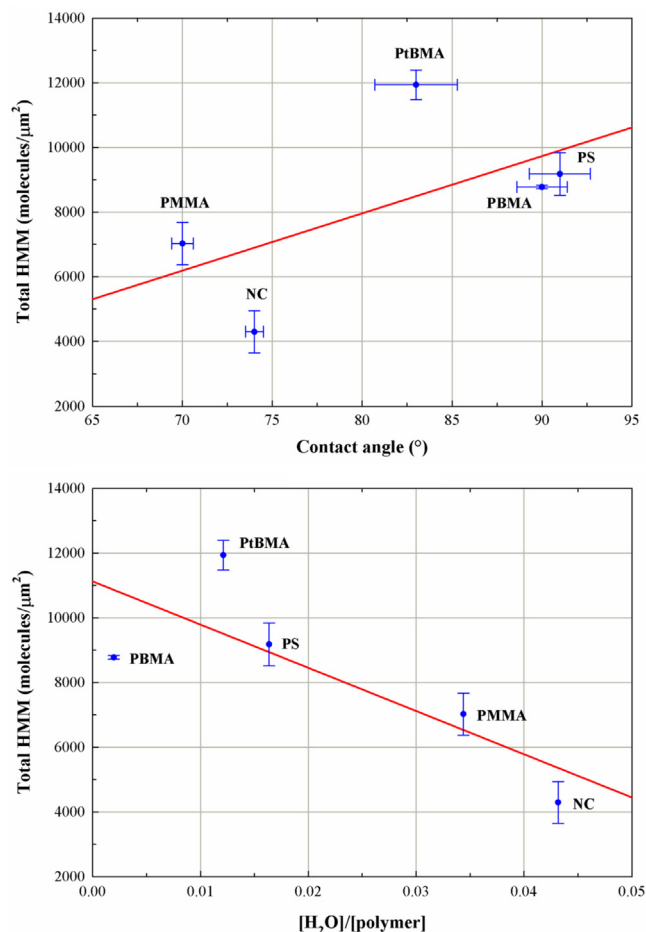


Fig. 2. Amount of HMM molecules adsorbed on various polymers as a measure of their surface hydrophobicity (a, top); and versus their water intake (b, bottom).

estimated using ATPase assay, has been reported only once (Albet-Torres et al., 2007a), and on non-polymeric surfaces. Interestingly, it has been observed that the average velocity of actin filaments increases with the hydrophobicity of the surfaces that immobilise myosin, if this surface is flat and rigid, e.g., silanised glass (Sundberg et al., 2006). However, in a clear sign that other factors are at play in determining the bioactivity of the motors, the average velocity of actin filaments on motor-functionalised polymers decreases with the increase of the surface hydrophobicity beyond a certain threshold (Nicolau et al., 2007b). Taking into account this contradictory information, the goal of the present study is to evaluate the impact of both polymer hydrophobicity and its viscoelastic properties on myosin density, activity, and motility of the actin filaments, to be able to appropriately select polymeric materials for molecular motors-based nanodevices.

3.2. Polymer properties control the amount of adsorbed HMM

Several studies (Fawcett et al., 1998; Hook et al., 2001, 2002; Rickert et al., 1997) demonstrated that protein layers adsorbed on surface incorporate substantial amounts of water. For instance, QCM measurements have demonstrated (Caruso et al., 1995; Geddes et al., 1994; Hook et al., 2002; Rickert et al., 1997) that, due to the additional mass of incorporated water, the frequency change following the protein adsorption in a liquid environment is 4x (Caruso et al., 1995; Geddes et al., 1994; Hook et al., 2002; Rickert et al., 1997) to 10x (Albet-Torres et al., 2007a; Hook et al., 2001) times higher than those in air. The QCM measurements of

adsorbed masses on rigid layers is nearly trivial, i.e., fully described by the classical Sauerbrey equation (Sauerbrey, 1959), but data interpretation for non-rigid layers is considerably more difficult because of the coupling of viscoelastic effects for soft, gel-like layers. However, the capacity of QCM methodology to measure both the adsorbed mass, as well as the viscoelastic properties of the protein layers, has been used to estimate these properties separately, e.g., the amount of motor protein on rigid surfaces (Albet-Torres et al., 2007a; Hanson et al., 2006), and the conformation of the protein motor layer as transducing the motility state in non-optical signal (Van Zalinge et al., 2012), respectively. Therefore, the measurements of (i) the absorbed water in the polymer layers, followed by (ii) the measurement of the adsorbed motors on the polymer surfaces, and finally (iii) the measurements of the absorbed water in the protein layers, is likely to progress the understanding of the behaviour of motor proteins immobilized on, or into polymeric layers.

The water uptake in the polymer film, i.e., before immobilisation of HMM, is generally lower for more hydrophobic polymers (Supplementary information, Fig. S11). While this relationship is generally expected, NC absorbs an unusual high amount of water, i.e., nearly 30% more than the next water-absorbing polymer (PMMA), with which it clusters as higher water-adsorbing polymers. This suggests that NC, and to a lesser extent PMMA, form thick gel-like layers in/on which the HMM molecules are immobilized. Equally expected, the water uptake is lower for polymers having higher glass transition (T_g), with the notable exception of PBMA (Supplementary information, Fig. S12). For glassy polymers, i.e., all tested less PBMA, a high T_g is the result of a higher stiffness of the polymer chains, which will decrease the rate of water diffusion. Conversely, for the rubbery PBMA, the much more flexible polymer network amplifies the impact of its hydrophobicity, thus leading to far less water uptake than it would be otherwise expected for its contact angle. Finally, the water uptake appears to increase, for glassy polymers, with their density (Supplementary information, Fig. S13). This apparently counter-intuitive trend can be understood by the larger content of hydrophilic groups, i.e., hydroxyl groups for NC, carboxyl groups for PMMA, which will result in stronger hydrogen bonding, and subsequently more compact polymeric networks and higher density.

A higher polymer surface hydrophobicity results in a higher amount of total adsorbed HMM (Fig. 2a). This strong relationship between surface hydrophobicity and the mass of adsorbed protein has been explained before in general terms (Ostuni et al., 2003; Vasina et al., 2009), as well as specifically regarding HMM adsorption (Albet-Torres et al., 2007a; Nicolau et al., 2007b; Sundberg et al., 2003; Van Zalinge et al., 2012). Interestingly, NC presents the lowest amount of adsorbed HMM, despite the fact that it is the standard surface used in motility assays studies. Conversely, PS, which has not reported as being used in motility studies, although it is commonly used in assays using surface-immobilized proteins, presents the highest adsorbed amount of HMM. Importantly for the understanding of the control of motility characteristics by polymer properties, a higher water uptake in the polymer, which is largely in an inverse relationship with the polymer hydrophobicity, albeit not in an univoque manner, results in a lower amount of HMM adsorbed on, or rather absorbed in, the polymer layer (Fig. 2b). Finally, the uptake of water in HMM layer correlates well with the water uptake in the polymer layer (Supplementary Information, Fig. S14).

3.3. Structure of the polymer-HMM layer

The estimation of the mass of surface-adsorbed HMM allows the calculation of the total surface density of HMM molecules (Table 1). The overall contribution to the acoustic impedance is

Table 1

Surfaces, total and active HMM surface density and actin filament motility parameters.

Surface	Contact angle (°)	Total HMM density ^a (molecules μm^{-2})	Active HMM density ^b (molecules μm^{-2})	[Active]/[total] HMM	Motile filaments (%)	Average velocity ($\mu\text{m s}^{-1}$)
PMMA	70 ± 0.6	7024 ± 652	2553 ± 296	0.36 ± 0.07	57	2.62 ± 1.22
NC	74 ± 0.5	4294 ± 645	2928 ± 1014	0.68 ± 0.16	75	3.45 ± 1.28
PtBMA	83 ± 2.3	11,935 ± 459	2442 ± 466	0.20 ± 0.03	48	1.68 ± 0.86
PBMA	90 ± 1.4	8773 ± 55	1312 ± 208	0.15 ± 0.02	55	2.20 ± 1.06
PS	91 ± 1.7	9177 ± 661	1046 ± 181	0.11 ± 0.03	61	2.15 ± 1.05
NC ^c	–	–	2509 ± 328	–	–	–
NC ^d	–	–	2400	–	–	–
TMCS ^e	75.1 ± 1.3	~ 20,000	5000–7000	~ 0.3	–	–

Notes.

Due to the large amount of water trapped in the HMM layer, which can be up to 10x the mass of the protein, the total number of HMM is a relative approximation. Additionally, TMCS is a very flat, very rigid surface, with low similarity with the polymer materials discussed here.

^a Calculated from QCM measurements, corrected for polymer effects.

^b Calculated from ATPase data.

^c Data from (Guo and Guilford, 2004).

^d Data from (Uyeda et al., 1991).

^e Data from (Albet-Torres et al., 2007a).

calculated by matrix multiplication, starting from the top layer, the HMM solution. The results of these calculations are presented in the [Supplementary information section \(Table SI2\)](#). One observation is that the density of the protein layer for all surfaces is approximately the same, i.e., $\sim 1090 \text{ kg m}^{-3}$, which is close to the values reported elsewhere (Hook et al., 2002), e.g., 1040 kg m^{-3} . However, there are notable differences in the thickness of the protein layers, which suggests different conformations of the motor protein on different surfaces, and possibly its surface-bound site. The values obtained for the total mass adsorbed and the number of heads, i.e., $280\text{--}540 \text{ ng/cm}^2$, which depend on the surface properties (higher on hydrophobic surfaces), are similar with those reported elsewhere (Harada, 1990), i.e., 350 ng/cm^2 for myosin on glass.

A vertically-structured, multi-layer model, comprising, from bottom up, [solid substrate]-[stiff polymer]-[polymer gel]-[HMM layer], has been recently proposed (Van Zalinge et al., 2012), but the present experiments allow a more precise description, as follows:

- (1) At the bottom of the multilayer, the basal surface supports polymer layer of $500\text{--}750 \text{ nm}$, which contains various amounts of water, as described in the previous sections.
- (2) This polymer layer supports a HMM-rich layer, which also contains water. This HMM layer comprises two sub-layers with different densities and dynamical moduli, thus

contributing differently to the acoustic impedance. A bottom sub-layer, several nm's thick, which due to its reduced thickness can be treated within the Sauerbrey approximation, is covered by a top sub-layer with a thickness of $\sim 15\text{--}35 \text{ nm}$, rich in water, comprising one or both S1 heads, and which has the characteristics of a gel, thus modelled as a viscoelastic layer.

- (3) The constitutive element of the polymer-protein layer, i.e., HMM molecule, is $80\text{--}90 \text{ nm}$ long, including the S1 heads, which are $\sim 19 \text{ nm}$ long with a diameter of $2\text{--}3 \text{ nm}$. These geometrical considerations, coupled with the large amount of water trapped in the protein layer suggest an arrangement of HMM molecules with the head groups tightly packed and the long axes of the motors perpendicular to the surface, as proposed before (Albet-Torres et al., 2007a).

The [solid substrate]-[stiff polymer]-[polymer gel]-[HMM layer] model described above is schematically presented in [Fig. 3](#).

3.4. Polymer properties control the amount of active HMM

The immediate observation when comparing the density of QCM-measured, total HMM molecules ([Fig. 2a](#)) and ATPase-measured, active HMM molecules ([Fig. 4a](#)), is that the surface hydrophobicity has a contradictory effect on these two parameters. Indeed, as proposed above and elsewhere (Nicolau et al., 2007b), the

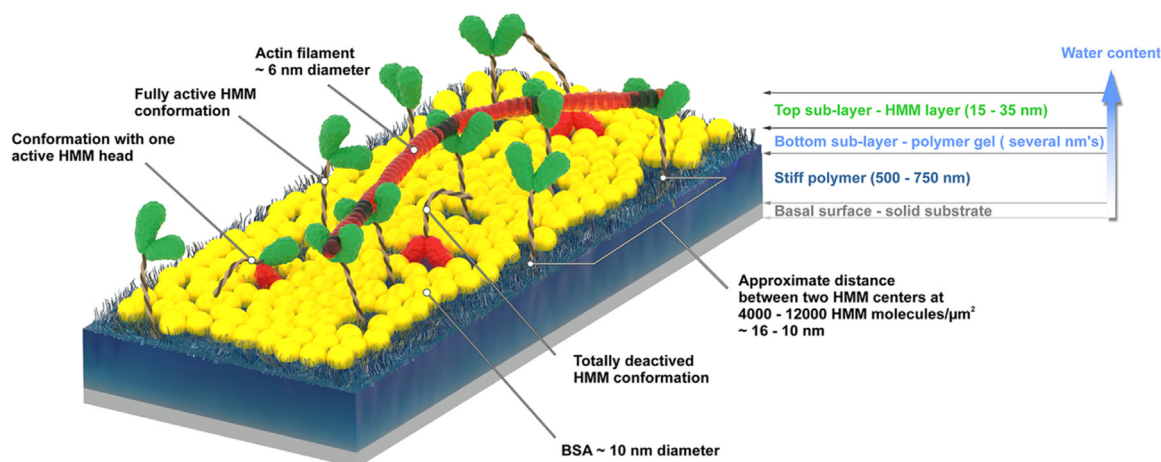


Fig. 3. Multi-layer model of the [solid substrate]-[stiff polymer]-[polymer gel]-[HMM layer].

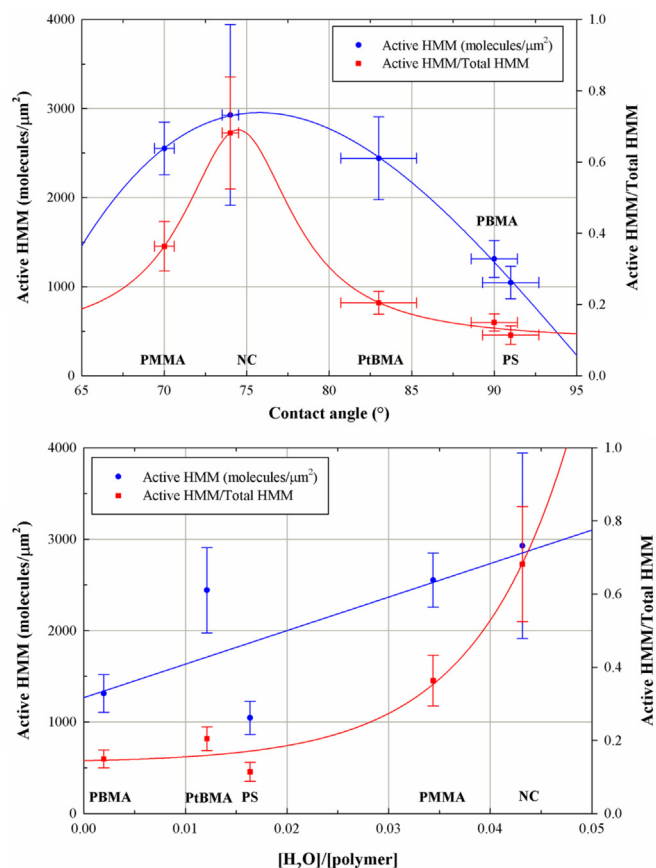


Fig. 4. Active HMM molecules adsorbed on various polymers, and the ratio of active/total HMM molecules, versus the respective surface hydrophobicities (a, top); and versus the water intake in the respective polymer (b, bottom).

surfaces with higher hydrophobicity induce the adsorption of a higher *total* number of HMM molecules, but also a lower number of *active* HMM molecules. This is not surprising, as hydrophobic surfaces will almost always induce a larger adsorption of proteins (Ostuni et al., 2003; Vasina et al., 2009). Moreover, the alternative-to-hydrophobicity driving force for protein adsorption, i.e., electrostatic interactions, will be more prevalent on hard, molecularly flat surfaces, e.g., silicon-based materials (Albet-Torres et al., 2007a), as the short-range hydrophobic interactions will be prevalent due to the brush-like nature of polymeric surfaces. Also, the hydrophobic interactions between the polymeric surfaces and motor proteins will result in the expulsion of water starting from the hydrophobic core outwards the protein, thus inducing its denaturation (Ostuni et al., 2003). The overall result of this 'tug-of-war' between larger *total* HMM molecules, but also lower *active* HMM molecules, is a decrease of the ratio of active vs. total HMM molecular density with surface hydrophobicity. An increase of the water uptake in the polymer (in negative correlation with polymer hydrophobicity), results in an increase of both active HMM molecules, and importantly the ratio between the active and total number of HMM molecules (Fig. 4b).

3.5. Polymer properties control the motility of actin filaments

The various amounts of total HMM molecules adsorbed on various polymer surfaces, and in particular the impact of the properties of respective polymers on the preservation of bioactivity of motor proteins, result in variations of the characteristics of the motility of actin filaments sliding on HMM-functionalised surfaces (Fig. 5, also Table 1). The analysis of the characteristics of

the motility of actin filaments on surfaces shows that only NC induces a rapid and smooth movement, indicated by the mono-modal distribution of velocities, skewed towards higher values, and monotonous distribution of deflection angles skewed towards 0 degrees. All other surfaces present bimodal distributions of velocities, which could be related to denaturated HMM stalling the motility, or a population of partially-denaturated HMM working with a single arm, which will propel actin filaments at lower velocities (Toyoshima et al., 1987). Also all surfaces other than NC (and possibly PMMA) induce a 'stuttered' motion of the actin filaments as evidenced by the broader band of frequency on both sides of 0 values for acceleration (an acceleration means perfectly smooth movement), and large proportion of non-zero deflection angles. This behaviour can be understood by the higher proportion of non-active HMM molecules, which will stall or brake the motion of actin filaments.

A clear, but initially counterintuitive trend, that is, of the decreasing velocity of actin filaments with the increase of the total HMM molecular density can be observed in Fig. 6a. Rather expectedly, the increase of the active HMM leads to an increase of the actin filament velocity (Fig. 6b). Finally, a stronger positive correlation exists between filament velocity and the proportion of active to total heads (Fig. 6c). Of particular note in these results is that NC, which is typically considered the 'gold standard' for actomyosin motility, presents by far the highest proportion of active per total HMM (also Table 1). The relationship in Fig. 6c is positive, consistent with previously described models (Cuda et al., 1997; Tawada and Sekimoto, 1991; Warshaw et al., 1990), which assume that the *in vitro* actomyosin velocity is the result of a balance between myosin cross-bridges in the productive force-generating state and external load imposed by non-productive cyclical actin-myosin interactions.

3.6. Selection of materials for dynamic nanodevices using molecular motors

In the view of the complex relationship between the properties of the surfaces immobilising protein molecular motors, in particular myosin, as described above, the general guidelines for the selection of the polymeric materials to be employed in the manufacturing of dynamic nanodevices can be summarised as follows:

- **Key performance criteria for motility.** Different dynamic nanodevices using protein molecular motors, would require different key performance features, depending on their intended application. For instance, diagnostic devices can capitalise on the much higher rate of interaction between the target and the probe molecules than that driven by diffusion, as in classical assays, e.g., ELISA (Nitta and Hess, 2005). For these devices, which transduce the biomolecular recognition into change of motility (Korten et al., 2010; Kumar et al., 2016; Martinez-Neira et al., 2005; Van Zalinge et al., 2012), and which will interact with various biological fluids (Korten et al., 2013), the maintenance of motility is paramount. On the other hand, in other applications, such as diagnostic devices using molecular motors to transport and concentrate the target molecules in selected locations (Kumar et al., 2016; Lard et al., 2013); or biocomputation devices (Nicolau et al., 2016), a high velocity of the cytoskeletal filaments could become more important. As a general rule, however, the maintenance of the motility, ensured by the presence of a large number of active motors, and/or the preservation of their bioactivity is critical for any successful implementation of dynamic nanodevices using protein molecular motors.
- **Surface hydrophobicity.** Materials with a high surface hydrophobicity will induce a high total amount of adsorbed motors,

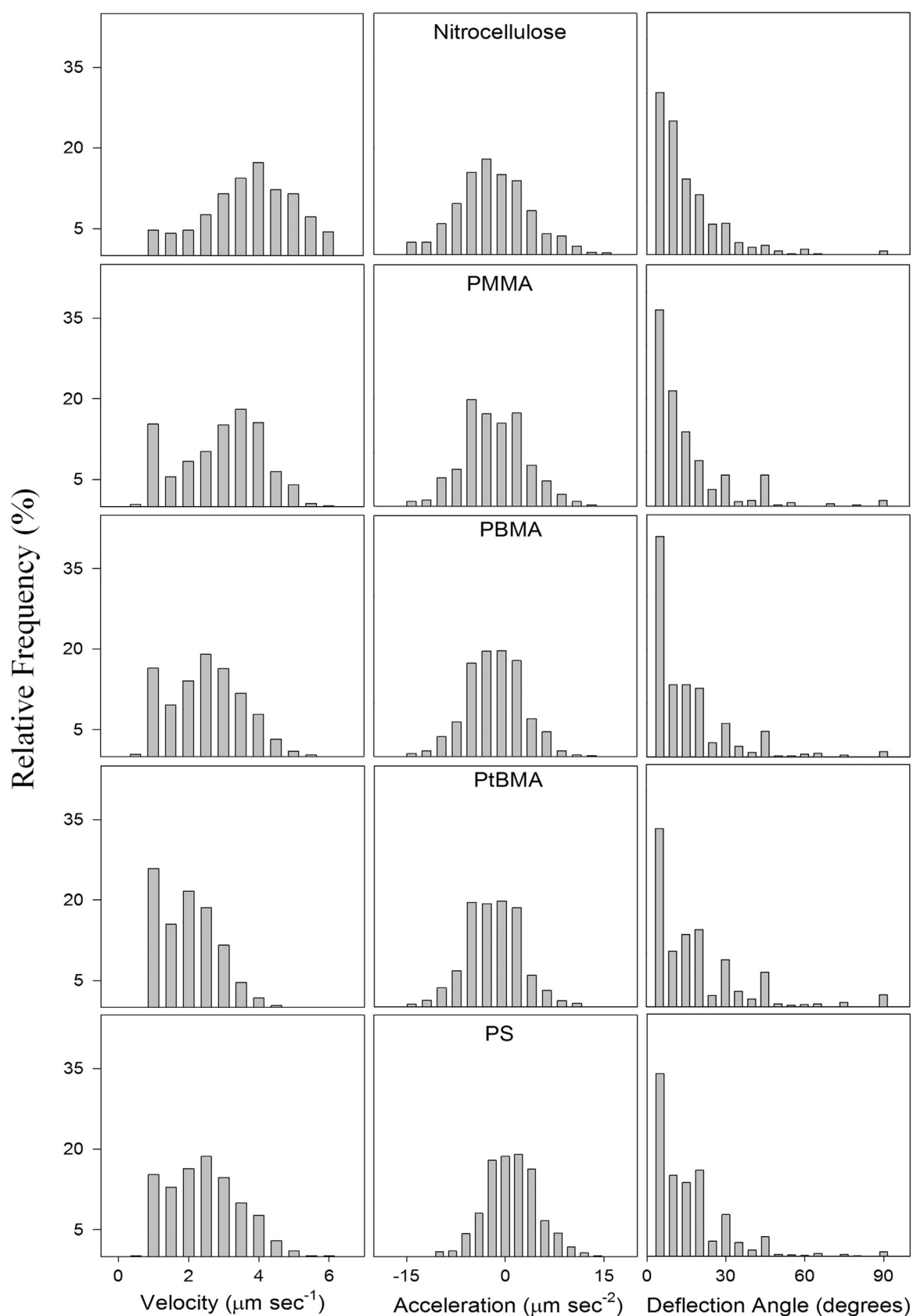


Fig. 5. Histograms of the frequency of co-occurrences of velocity, velocity variation; and deflection angle of the motility of actin filaments on various HMM-functionalised polymeric surfaces.

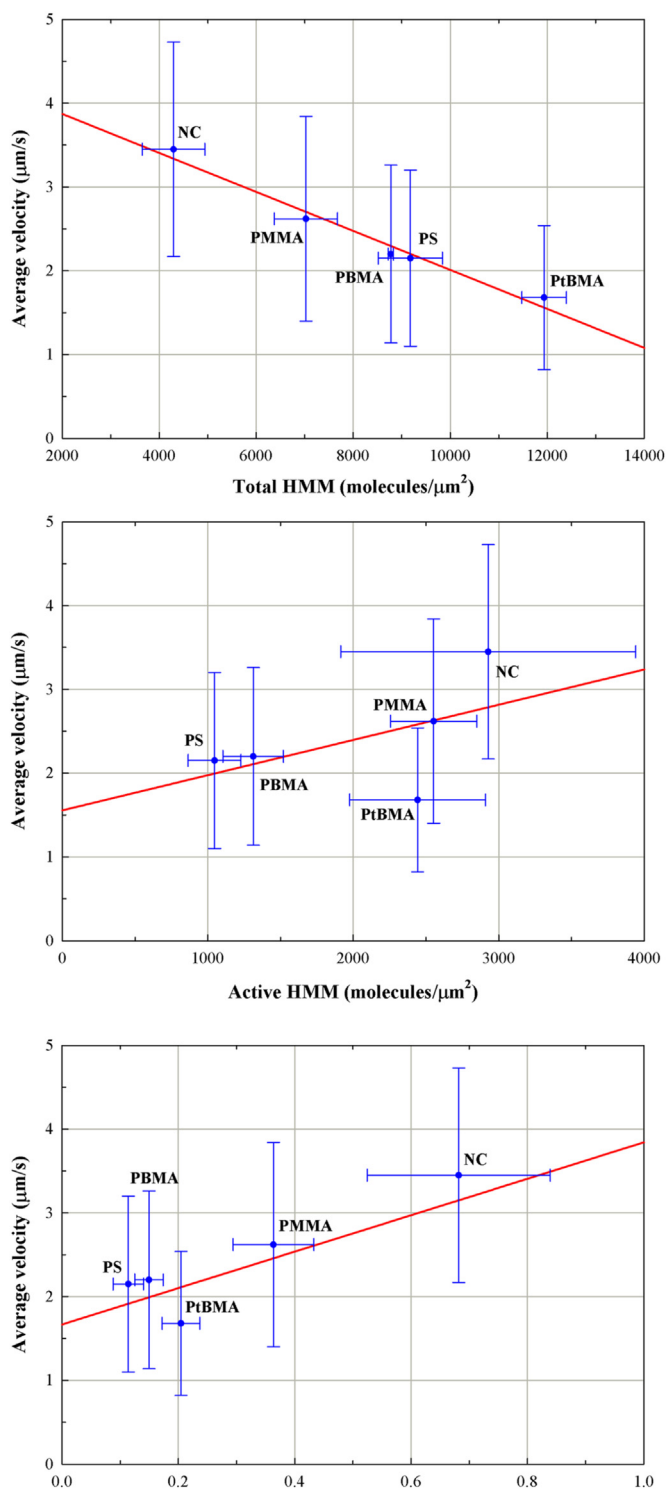


Fig. 6. Relationship between the average velocity of the actin filaments and total number of HMM molecules (a, top); active number of HMM molecules (b, middle); and the ratio of active per total number of HMM (c, bottom).

which would be expected to translate in a higher density of active motors. However, for polymeric materials, it was demonstrated here that a higher hydrophobicity translates in high rates of denaturation, which greatly overshadow the benefits of higher overall total density of motors. Furthermore, denaturated motors could have a deleterious effect on motility. Consequently, for polymeric surfaces, a decreased hydrophobicity is likely to be beneficial for the motility of cytoskeletal filaments.

It appears that this inverse relationship between surface hydrophobicity and preservation of the bioactivity of the motors is specific to polymeric materials, as it has been demonstrated that of very flat, very rigid surfaces, e.g., silanised silicon oxide (Albet-Torres et al., 2007a), a higher hydrophobicity results in higher amounts of total and active motors. While these rigid surfaces appear to address the key performance criteria regarding motility, polymeric materials offer important advantages regarding fabrication, cost and disposability (Becker and Gärtner, 2000; Sollier et al., 2011).

- **Water uptake in polymers.** It appears that the preservation of the activity of the motors is not only due to a low hydrophobicity of the material, but also by the capacity of the polymer to form gel-like films via water uptake. Indeed, while NC is not the most hydrophilic polymer studied here, its very high capacity for water uptake, even comparing with the nearest candidate (PMMA) makes it the best material regarding the preservation of motor activity.
- **New materials tailored for dynamic nanodevices.** The present study demonstrated the physico-chemical reasons why NC is the 'gold standard' of motility assays, i.e., moderate hydrophobicity, and high level of water uptake. However, for microfluidics devices, NC is far from being an optimal material for their fabrication and operation, e.g., due to its inflammability. The key performance parameters demonstrated here, i.e., moderate hydrophobicity, and higher capacity for water uptake, as well as other, e.g., optical transparency, could be the leads for the search of other materials for dynamic nanodevices. For instance, cyclo-olefin copolymers (COC) have been demonstrated to be appropriate for micro-manufacturing (Mair et al., 2006), and also to have the appropriate physico-chemical properties for biomolecule immobilisation (Nicolau et al., 2005).

4. Conclusions

We have investigated the density and bioactivity of HMM on a variety of polymer surfaces, and related the results to in vitro actomyosin motility velocities. The results show that the proportion of active per total HMM heads, rather than the surface density of active HMM, controls the actin filament sliding velocities. We have thus shown that the mechanism by which surfaces modulate actomyosin motility in vitro is complex, such that HMM molecules adsorbed to a surface in a perturbed conformation or ineffective orientation appear to play a significant role in decreasing filament sliding velocities. The results have application to the fundamental understanding of actin and myosin interactions, and to the future engineering of surfaces for molecular motor based nano-devices.

Contributions of the authors

KLH run the motility and ATPase assays. GS collected the motility data. SD run the QCM experiments. KLH, GS, FF, SD, OK, VT and DVN analysed the data. KLH, FF, OK, and DVN wrote the paper.

Acknowledgements

This work has been supported by grants from the Defence Advanced Research Projects Agency, under Grant agreements N66001-00-1-8952 and N66001-03-1-8913; the Australian Research Council (Discovery project DP0208754); European Union Seventh Framework Programme (FP7/2007–2011) under Grant agreement number 228971 (MONAD); and Natural Science and

Engineering Research Council (NSERC, Canada), Grant RGPIN-2016-05019. The authors wish to thank Dr. Murat Kekic and Prof. Cristobal dos Remedios from the University of Sydney, and Dr. Jenny Aveyard from the University of Liverpool, for myosin preparation; and to Dr. Harm van Zalinge for assistance with QCM measurements.

Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at doi:10.1016/j.bios.2016.08.061.

References

- Agarwal, A., Katira, P., Hess, H., 2009. Millisecond curing time of a molecular adhesive causes velocity-dependent cargo-loading of molecular shuttles. *Nano Lett.* 9 (3), 1170–1175.
- Albet-Torres, N., Gunnarsson, A., Persson, M., Balaz, M., Höök, F., Månsson, A., 2010. Molecular motors on lipid bilayers and silicon dioxide: different driving forces for adsorption. *Soft Matter* 6 (14), 3211–3219.
- Albet-Torres, N., O'Mahony, J., Charlton, C., Balaz, M., Lisboa, P., Aastrup, T., Månsson, A., Nicholls, I.A., 2007a. Mode of heavy meromyosin adsorption and motor function correlated with surface hydrophobicity and charge. *Langmuir* 23, 11147–11156.
- Albet-Torres, N., O'Mahony, J., Charlton, C., Balaz, M., Lisboa, P., Aastrup, T., Månsson, A., Nicholls, I.A., 2007b. Mode of heavy meromyosin adsorption and motor function correlated with surface hydrophobicity and charge. *Langmuir* 23 (22), 11147–11156.
- Bakewell, D.J.G., Nicolau, D.V., 2007. Protein linear molecular motor-powered nanodevices. *Aust. J. Chem.* 60 (5), 314–332.
- Barden, J.A., Dosremedios, C.G., 1984. The environment of the high-affinity cation binding site on actin and the separation between cation and ATP sites as revealed by protein NMR and fluorescence spectroscopy. *J. Biochem.* 96 (3), 913–921.
- Becker, H., Gärtner, C., 2000. Polymer microfabrication methods for microfluidic analytical applications. *Electrophoresis* 21 (1), 12–26.
- Bull, J.L., Hunt, A.J., Meyhofer, E., 2005. A theoretical model of a molecular-motor-powered pump. *Biomed. Microdevices* 7 (1), 21–33.
- Bunk, R., Klinth, J., Montelius, L., Nicholls, I.A., Omling, P., Tagerud, S., Månsson, A., 2003a. Actomyosin motility on nanostructured surfaces. *Biochem. Biophys. Res. Commun.* 301 (3), 783–788.
- Bunk, R., Sundberg, M., Månsson, A., Nicholls, I.A., Omling, P., Tagerud, S., Montelius, L., 2005. Guiding motor-propelled molecules with nanoscale precision through silanized bi-channel structures. *Nanotechnology* 16 (6), 710–717.
- Bunk, R., Klinth, J., Rosengren, J., Nicholls, I., Tagerud, S., Omling, P., Månsson, A., Montelius, L., 2003b. Towards a 'nano-traffic' system powered by molecular motors. *Microelectron. Eng.* 67–8, 899–904.
- Carsten, M.E., Mommaerts, W.F.H., 1963. A study of actin by means of starch gel electrophoresis. *Biochemistry* 2 (1), 28–8.
- Caruso, F., Serizawa, T., Furlong, D.N., Okahata, Y., 1995. Quartz-crystal microbalance and surface-plasmon resonance study of surfactant adsorption onto gold and chromium-oxide surfaces. *Langmuir* 11 (5), 1546–1552.
- Clemmens, J., Hess, H., Howard, J., Vogel, V., 2003a. Analysis of microtubule guidance in open microfabricated channels coated with the motor protein kinesin. *Langmuir* 19 (5), 1738–1744.
- Clemmens, J., Hess, H., Doot, R., Matzke, C.M., Bachand, G.D., Vogel, V., 2004. Motor-protein "roundabouts": microtubules moving on kinesin-coated tracks through engineered networks. *Lab Chip* 4 (2), 83–86.
- Clemmens, J., Hess, H., Lipscomb, R., Hanein, Y., Bohringer, K.F., Matzke, C.M., Bachand, G.D., Bunker, B.C., Vogel, V., 2003b. Mechanisms of microtubule guiding on microfabricated kinesin-coated surfaces: chemical and topographic surface patterns. *Langmuir* 19 (26), 10967–10974.
- Cuda, G., Pate, E., Cooke, R., Sellers, J.R., 1997. In vitro actin filament sliding velocities produced by mixtures of different types of myosin. *Biophys. J.* 72 (4), 1767–1779.
- Fawcett, N.C., Craven, R.D., Zhang, P., Evans, J.A., 1998. QCM Response to Solvated, Tethered Macromolecules. *Anal. Chem.* 70, 2876–2880.
- Fischer, T., Agarwal, A., Hess, H., 2009. A smart dust biosensor powered by kinesin motors. *Nat. Nanotechnol.* 4 (3), 162–166.
- Fulga, F., Nicolau, J., Nicolau, D.V., 2009. Models of protein linear molecular motors for dynamic nanodevices. *Integr. Biol.* 1 (2), 150–169.
- Geddes, N.J., Paschinger, E.M., Furlong, D.N., Ebara, Y., Okahata, Y., Than, K.A., Edgar, J.A., 1994. Piezoelectric crystal for the detection of immunoreactions in buffer solutions. *Sens. Actuators B Chem.* 17 (2), 125–131.
- Guo, B., Guilford, W.H., 2004. The tail of myosin reduces actin filament velocity in the in vitro motility assay. *Cell Motil. Cytoskeleton* 59 (4), 264–272.
- Hanson, K.L., Viidyanathan, V., Nicolau, D.V., 2006. Actomyosin Motility Detection using Quartz Crystal Microbalance.
- Harada, Y., 1990. Mechanochemical coupling in actomyosin energy transduction studied by in vitro movement assay. *J. Mol. Biol.* 216, 49–68.
- Harris, D.E., Warshaw, D.M., 1993. Smooth and skeletal-muscle myosin both exhibit low duty cycles at zero load in-vitro. *J. Biol. Chem.* 268 (20), 14764–14768.
- Hiratsuka, Y., Tada, T., Oiwa, K., Kanayama, T., Uyeda, T.Q.P., 2001. Controlling the direction of kinesin-driven microtubule movements along microlithographic tracks. *Biophys. J.* 81 (3), 1555–1561.
- Holzbaur, E.L., Goldman, Y.E., 2010. Coordination of molecular motors: from in vitro assays to intracellular dynamics. *Curr. Opin. Cell Biol.* 22 (1), 4–13.
- Hook, F., Kasemo, B., Nylander, T., Fant, C., Sott, K., Elwing, H., 2001. Variations in coupled water, viscoelastic properties, and film thickness of a Mefp-1 protein film during adsorption and cross-linking: a quartz crystal microbalance with dissipation monitoring, ellipsometry, and surface plasmon resonance study. *Anal. Chem.* 73 (24), 5796–5804.
- Hook, F., Voros, J., Rodahl, M., Kurrat, R., Boni, P., Ramsden, J.J., Textor, M., Spencer, N.D., Tengvall, P., Gold, J., Kasemo, B., 2002. A comparative study of protein adsorption on titanium oxide surfaces using in situ ellipsometry, optical waveguide lightmode spectroscopy, and quartz crystal microbalance/dissipation. *Colloids Surf. B-Biointerfaces* 24 (2), 155–170.
- Huxley, A.F., Niedergerke, R., 1954. Structural changes in muscle during contraction - Interference microscopy of living muscle fibres. *Nature* 173 (4412), 971–973.
- Huxley, H., Hanson, J., 1954. Changes in the cross-striations of muscle during contraction and stretch and their structural interpretation. *Nature* 173 (4412), 973–976.
- Jaber, J.A., Chase, P.B., Schlenoff, J.B., 2003. Actomyosin-driven motility on patterned polyelectrolyte mono- and multilayers. *Nano Lett.* 3 (11), 1505–1509.
- Katira, P., Agarwal, A., Hess, H., 2009. A random sequential adsorption model for protein adsorption to surfaces functionalized with poly(ethylene oxide). *Adv. Mater.* 21 (16), 1599–+.
- Katira, P., Agarwal, A., Fischer, T., Chen, H.Y., Jiang, X., Lahann, J., Hess, H., 2007. Quantifying the performance of protein-resisting surfaces at ultra-low protein coverages using kinesin motor proteins as probes. *Adv. Mater.* 19 (20), 3171–3176.
- Kinbara, K., Aida, T., 2005. Toward intelligent molecular machines: directed motions of biological and artificial molecules and assemblies. *Chem. Rev.* 105 (4), 1377–1400.
- Korten, S., Albet-Torres, N., Paderi, F., Ten Siethoff, L., Diez, S., Korten, T., Te Kronnie, G., Månsson, A., 2013. Sample solution constraints on motor-driven diagnostic nanodevices. *Lab Chip Miniat. Chem. Biol.* 13 (5), 866–876.
- Korten, T., Månsson, A., Diez, S., 2010. Towards the application of cytoskeletal motor proteins in molecular detection and diagnostic devices. *Curr. Opin. Biotechnol.* 21 (4), 477–488.
- Kron, S.J., Spudich, J.A., 1986. Fluorescent actin filaments move on myosin fixated to a glass surface. *Proc. Natl. Acad. Sci. USA* 83 (17), 6272–6276.
- Kumar, S., Milani, G., Takatsuki, H., Lana, T., Persson, M., Frasson, C., Te Kronnie, G., Månsson, A., 2016. Sensing protein antigen and microvesicle analytes using high-capacity biopolymer nano-carriers. *Analyst* 141 (3), 836–846.
- Lard, M., ten Siethoff, L., Kumar, S., Persson, M., te Kronnie, G., Linke, H., Månsson, A., 2013. Ultrafast molecular motor driven nanoseparation and biosensing. *Biosens. Bioelectron.* 48, 145–152.
- Limberis, L., Stewart, R.J., 2000. Toward kinesin-powered microdevices. *Nanotechnology* 11 (2), 47–51.
- Mair, D.A., Geiger, E., Pisano, A.P., Fréchet, J.M.J., Svec, F., 2006. Injection molded microfluidic chips featuring integrated interconnects. *Lab Chip Miniat. Chem. Biol.* 6 (10), 1346–1354.
- Margossian, S.S., Lowey, S., 1982. Preparation of myosin and its subfragments from rabbit skeletal-muscle. *Methods Enzymol.* 85, 55–71.
- Martinez-Neira, R., Kekic, M., Nicolau, D., dos Remedios, C.G., 2005. A novel biosensor for mercuric ions based on motor proteins. *Biosens. Bioelectron.* 20 (7), 1428–1432.
- Moorjani, S.G., Jia, L., Jackson, T.N., Hancock, W.O., 2003. Lithographically patterned channels spatially segregate kinesin motor activity and effectively guide microtubule movements. *Nano Lett.* 3 (5), 633–637.
- Nicolau, D.V., Suzuki, H., Mashiko, S., Taguchi, T., Yoshikawa, S., 1999. Actin motion on microlithographically functionalized myosin surfaces and tracks. *Biophys. J.* 77 (2), 1126–1134.
- Nicolau, D.V., Pham, D.K., Ivanova, E.P., Wright, J.P., Lenigk, R., Smekal, T., Grodzinski, P., 2005. Tone reversal of an AFM lateral force image due to hybridization of oligonucleotides immobilized on polymers. *Small* 1 (6), 610–613.
- Nicolau, D.V., Solana, G., Kekic, M., Fulga, F., Mahanivong, C., Wright, J., dos Remedios, C.G., 2007a. Surface Hydrophobicity modulates the operation of actomyosin-based dynamic nanodevices. *Langmuir* 23, 10846–10854.
- Nicolau, D.V., Solana, G., Kekic, M., Fulga, F., Mahanivong, C., Wright, J., Ivanova, E.P., dos Remedios, C.G., 2007b. Surface hydrophobicity modulates the operation of actomyosin-based dynamic nanodevices. *Langmuir* 23, 10846–10854.
- Nicolau, J., Lard, D.V., Korten, M., Van Delft, T., Persson, F.C.M.J.M., Bengtsson, M., Månsson, E., Diez, A., Linke, S., Nicolau, D.V., H., 2016. Parallel computation with molecular-motor-propelled agents in nanofabricated networks. *Proc. Natl. Acad. Sci. USA* 113 (10), 2591–2596.
- Nitta, T., Hess, H., 2005. Dispersion in active transport by kinesin-powered molecular shuttles. *Nano Lett.* 5 (7), 1337–1342.
- Ostuni, E., Grzybowski, B.A., Mrksich, M., Roberts, C.S., Whitesides, G.M., 2003. *Langmuir* 19, 1861–1872.
- Rickert, J., Brecht, A., Gopel, W., 1997. Quartz crystal microbalances for quantitative biosensing and characterizing protein multilayers. *Biosens. Bioelectron.* 12 (7), 567–575.
- Rivelino, D., Ott, A., Julicher, F., Winkelmann, D.A., Cardoso, O., Lacapere, J.J.,

- Magnusdottir, K.M., Viovy, J.L., Gorre-Talini, L., Prost, J., 1998. Acting on actin: the electric motility assay. *Eur. Biophys. J. Biophys. Lett.* 27 (4), 403–408.
- Sauerbrey, G., 1959. Verwendung von schwingquartzen zur wagung dunner schichten und zur mikrowagung. *Z. Phys.* 155 (2), 206–222.
- Sellers, J.R., Cuda, G., Wang, F., Homsher, E., 1993. Myosin-specific adaptations of the motility assay. *Methods Cell Biol.* 39 (39), 23–49.
- Sollier, E., Murray, C., Maoddi, P., Di Carlo, D., 2011. Rapid prototyping polymers for microfluidic devices and high pressure injections. *Lab Chip Miniat. Chem. Biol.* 11 (22), 3752–3765.
- Spudich, J.A., 2011. Molecular motors: forty years of interdisciplinary research. *Mol. Biol. Cell* 22 (21), 3936–3939.
- Spudich, J.A., Watt, S., 1971. Regulation of rabbit skeletal contraction 1. Biochemical studies on interaction of tropomyosin-troponin complex with actin and proteolytic fragments of myosin. *J. Biol. Chem.* 246 (15), 4866–8.
- Sundberg, M., Rosengren, J.P., Bunk, R., Lindahl, J., Nicholls, I.A., Tagerud, S., Omling, P., Montelius, L., Mansson, A., 2003. Silanized surfaces for in vitro studies of actomyosin function and nanotechnology applications. *Anal. Biochem.* 323 (1), 127–138.
- Sundberg, M., Balaz, M., Bunk, R., Rosengren-Holmberg, J.P., Montelius, L., Nicholls, I.A., Omling, P., Tagerud, S., Mansson, A., 2006. Selective spatial localization of actomyosin motor function by chemical surface patterning. *Langmuir* 22 (17), 7302–7312.
- Suzuki, H., Yamada, A., Oiwa, K., Nakayama, H., Mashiko, S., 1997. Control of actin moving trajectory by patterned poly(methyl methacrylate) tracks. *Biophys. J.* 72 (5), 1997–2001.
- Taussky, H.H., Shorr, E., 1953. A microcolorimetric method for the determination of inorganic phosphorus. *J. Biol. Chem.* 202 (2), 675–685.
- Tawada, K., Sekimoto, K., 1991. A physical model of atp-induced actin-myosin movement invitro. *Biophys. J.* 59 (2), 343–356.
- Toyoshima, Y.Y., Kron, S.J., McNally, E.M., Niebling, K.R., Toyoshima, C., Spudich, J.A., 1987. Myosin subfragment-1 is sufficient to move actin filaments in vitro. *Nature* 328 (6130), 536–539.
- Uyeda, T.Q.P., Kron, S.J., Spudich, J.A., 1990. Myosin step size - estimation from slow sliding movement of actin over low-densities of heavy-meromyosin. *J. Mol. Biol.* 214 (3), 699–710.
- Uyeda, T.Q.P., Warrick, H.M., Kron, S.J., Spudich, J.A., 1991. Quantized velocities at low myosin densities in an invitro motility assay. *Nature* 352 (6333), 307–311.
- Vale, R.D., 2003. The molecular motor toolbox for intracellular transport. *Cell* 112 (4), 467–480.
- Van Zalinge, H., Aveyard, J., Hajne, J., Persson, M., Mansson, A., Nicolau, D.V., 2012. Actin filament motility induced variation of resonance frequency and rigidity of polymer surfaces studied by quartz crystal microbalance. *Langmuir* 28 (42), 15033–15037.
- Vasina, E.N., Paszek, E., Nicolau, D.V., 2009. The BAD project: data mining, database and prediction of protein adsorption on surfaces. *Lab Chip* 9 (7), 891–900.
- Veigel, C., Schmidt, C.F., 2011. Moving into the cell: single-molecule studies of molecular motors in complex environments. *Nat. Rev. Mol. Cell Biol.* 12 (3), 163–176.
- Warshaw, D.M., Desrosiers, J.M., Work, S.S., Trybus, K.M., 1990. Smooth-muscle myosin cross-bridge interactions modulate actin filament sliding velocity in-vitro. *J. Cell Biol.* 111 (2), 453–463.

SUPPLEMENTARY INFORMATION

Polymer surface properties control the function of heavy meromyosin in dynamic nanodevices

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1. Polymers

The polymers on which HMM was immobilized for the testing of the in vitro motility of actin filaments covered a wide range of material properties, i.e., glass transition temperature (T_g); density; surface tension, surface free energy and contact angle (at equilibrium, advancing and receding), solubility index and shear stress moduli. These properties for the polymers selected are presented on SI Table 1.

All polymers with the exception of PBMA are in the glassy state at room temperature. The contact angle values in parentheses are those measured in this study. All other parameters, less the shear moduli are collected from an open source database (from http://www.accudynetest.com/polytable_01.html?sortby=sort_critical a references therein).

SI Table 1. Physical properties of the polymers and adsorbed HMM layers

Parameter	NC	PS	PMMA	PBMA	PtBMA
T _g [°C]	53	100	82	15	107
Polymer density [kg/m ³]	1060	1031	1060	1108	1020
Surface tension [mJ/m ²]	-	34.0	37.5	29.8	18.1
Surface free energy [mJ/m ²]	42.7	38.3	41.8	33.1	30.5
Contact angle, equilibrium	- (74)	87.4 (91)	69.1 (70)	91.0 (90)	108.1 (83)
Contact angle, advancing	54.7	88.5	74.7	-	-
Contact angle, receding	20.0	81.8	54.2	-	-
Hansen solubility index	23.5	23.9	21.3	18.4	-

2. Quartz Crystal Microbalance

2.1.Quartz Crystal Microbalance (QCM) equipment

The measurements of the amount of HMM adsorbed on the model surfaces used a commercial QCM system (QCM-Z500, from KSV Instruments), to measure the adsorbed protein mass and to assess more extensively the dynamics of viscoelastic polymers in thin films with thicknesses in the mesoscopic scale (hundreds of nanometers range). This advanced system allowed the

measurement of the impedance spectrum, thus providing both the frequency and the bandwidth, and addressing different, up to 11th harmonics.

The measurements have been performed sequentially, in a step-wise manner, i.e., first on the bare dry polymer surfaces, then on surfaces interfaced with the buffer solution and finally on surfaces with HMM adsorbed on the surfaces. A more elaborate description of the QCM measurement protocols and associated theoretical background is presented in the Supplementary Information Section.

2.2.QCM experimental protocols

The QCM measurement protocol consists of:

- (i) After mounting it in the probe, the bare crystal was allowed to stabilize and its frequency and resistance were recorded as control baseline. The operation was repeated several times to ensure that the baseline is not severely affected by the manipulation of the crystal. Furthermore, the crystals have been placed on the spin coater and mock spin coating tests have been run for the same purpose.
- (ii) The crystal was then carefully removed from the probe tested polymer was then spin-coated on the bare crystal, following the procedure described in the Methods section. Then the coated quartz crystal was mounted again in the probe, allowed to oscillate and its frequency and resistance oscillating in air was recorded. The polymer-coated crystal frequency and resistance were recorded as a second baseline.
- (iii) Buffer A was pumped into the probe chamber and the frequency and resistance shifts upon immersion were recorded as before.
- (iv) Finally, after stabilizing, HMM in Buffer A (0.1 mg mL⁻¹) was introduced in the flow cell. The crystal was allowed to equilibrate, then rinsed with Buffer A to remove reversibly the excess adsorbed protein, and the new shifts in frequency and resistance were recorded. The KSV-500 system recorded the temporal evolution of the impedance vs. frequency spectrum.

2.3.QCM theory

In air and for special properties of the deposited film, the relationship between the frequency change (Δf) and the amount of mass deposited on the quartz electrode is given by the Sauerbrey equation,(Sauerbrey 1959)

$$\Delta f = \frac{-2f_0^2 \Delta m}{A \sqrt{\rho_q \mu_q}} \quad (\text{Eq. SI1})$$

where f_0 is the fundamental frequency of the crystal (5×10^6 Hz for this study), Δm is the change in mass (g), A is the electrode area (cm²), ρ_q is the density of quartz (2.65 g cm⁻³) and μ_q is the shear modulus of quartz (2.95×10^{11} dyn cm⁻²). As it was shown(Banda et al. 2006; Benes 1984; Johannsmann 1999; Mason 1948) that the polymers behave like viscoelastic materials, the resonant frequency also depends on material properties(Granstaff and Martin 1994; Martin et al. 1994). In the case of non-rigid coatings, such as proteins, and even more for those also comprising polymers,

the added protein layer increases both the mass loading and the energy dissipation of the QCM,(Shen et al. 2001) and the increased energy dissipation will cause an additional change in frequency above that due to the mass loading alone.(Yang and Thompson 1993) Therefore, to convert the frequency shifts measured by the quartz crystal microbalance to deposited mass, the classical Sauerbrey equation (Sauerbrey 1959) valid for homogeneous, rigid, thin films, cannot be used. Moreover, it is known that on surfaces interfacing a liquid, as in our experiments, an additional frequency shift is observed due to liquid viscosity and density.(Kanazawa and Gordon 1985)

In order to make a quantitative assessment of the changes in the resonant frequency and in the motional resistance, an approximation to the transmission line model, which relates these changes to the acoustic load impedance Z_L acting on the surface of the crystal,(Lucklum et al. 1999, 2000; Lucklum and Hauptmann 2000a) has been used. From the acoustic load impedance, the following relations are derived:

$$\begin{aligned}\frac{\Delta f}{f_0} &= -\frac{\text{Im}(Z_L)}{\pi Z_q} \\ \frac{\Delta R}{2\omega L_q} &= \frac{\text{Re}(Z_L)}{\pi Z_q}\end{aligned}\quad (\text{Eq. SI2})$$

Here L_q is the motional inductance of the quartz crystal, specific for the used device; Z_q is the characteristic impedance of the crystal; and f is the resonance frequency of the bare quartz.

In a one-dimensional approximation, a layer with viscoelastic properties is completely characterized by its density ρ_1 , its thickness h_1 , as well as its shear storage modulus G' and its shear loss modulus G'' ($G_1 = G' + jG''$). The acoustic load of a viscoelastic film with finite thickness and without any other load at its surface is given by (Lucklum and Hauptmann 2000a)

$$Z_{L1} = jZ_{c1} \tan\left(\omega \sqrt{\frac{\rho_1}{G_1}} h_1\right); \quad Z_{c1} = \sqrt{\rho_1 G_1} \quad (\text{Eq. SI3})$$

where Z_{c1} is the characteristic impedance of the layer. Multilayer arrangements are treated with a matrix multiplication starting with the known acoustic load impedance at the top surface(Lucklum and Hauptmann 2000b). For two layers, the acoustic load is Z_L , the acoustic load impedance of the layer adjacent to the crystal.

$$Z_L = \frac{Z_{L1} + Z_{L2}}{1 + \frac{Z_{L1} \cdot Z_{L2}}{Z_{c1}^2}} \quad (\text{Eq. SI4})$$

2.4.QCM calculation protocols

First, the influence of the polymer layer in the QCM measurement was factored out. For each bare quartz crystal the resonance frequency and the motional resistance was measured in air; then the polymer was spin-coated on the quartz crystal according to the procedure described above and the

frequency shift and the motional resistance were measured. Following the procedure, buffer A was pumped into the chamber, and the new frequency shift and motional resistance were measured after the system stabilized. These measurements were performed for the 1st, 3rd, 5th, 7th and 11th harmonics.

The surface acoustic impedance was calculated as the contribution of the polymer layer for the first measurement performed in air, and of two layers, the polymer layer and the liquid, assumed infinitely thick, for the second case. Using the relations coupling the real and the imaginary part of the impedance to the frequency shift and the change in motional resistance, an optimization procedure was then used to estimate (i) the shear storage modulus G' ; (ii) the shear loss modulus G'' ; (iii) the density and thickness of the polymer films.

In the second step, HMM in Buffer A was pumped into the chamber and a new –protein- layer is formed between the polymer and the liquid. The impedance was then used in a new optimization procedure. Several parameters for each layer were used as reported in the literature. The shear modulus was assumed as 65000 N m^{-2} (Hook et al. 2001) and the viscosity was found as a result of the fitting procedure. Since the viscosity of 0.1 mg mL^{-1} HMM solution should not be significantly different from that of buffer alone, (Dainty et al. 1944) the top liquid layer was assumed to have the same parameters.

The results are presented in the Supplementary Information Table 1.

Supp. Info. Table 2. Physical properties of the polymers and adsorbed HMM layers

Parameter	PS	PMMA	PBMA	PtBMA	NC
Mass [ng/cm^2]	600	370	650	774	322
Mass gel [ng/cm^2]	176	78	171	235	41
Mass Sauerbrey [ng/cm^2]	424	292	480	539	281
Protein thickness [nm]	26.4	18.6	23.6	27.3	11.5
Protein density [kg/m^3]	1080	1070	1098	1100	1099
G' protein [N/m^2]	55000	55000	55000	55000	55000
eta protein [Ns/m^2]	0.0097	0.018	0.0061	0.014	0.0056
Polymer thickness [μm]	0.738	0.526	0.560	0.525	0.739
Polymer density [kg/m^3]	1031	1060	1108	1020	1060
G' [N/m^2], $\times 10^{-7}$	1.1	1.59	1.83	1.6	1.35
G'' [N/m^2] $\times 10^{-6}$	8.1	7.8	1.39	4.4	2.86
$\tan (G''/G')$	0.736	0.491	0.760	0.275	0.212
Tg [$^{\circ}\text{C}$]	100	82	15	107	53

2.5.QCM characterization of water uptake in polymers

Various polymers differ in terms of physical and chemical properties, as detailed in Table SI1, and this specificity translate in various response of the polymers with regard to water uptake, which in turn modulates the behavior versus HMM adsorption and preservation of its activity.

A synthetic view of the water uptake response for various polymers is presented in Figure SI1.

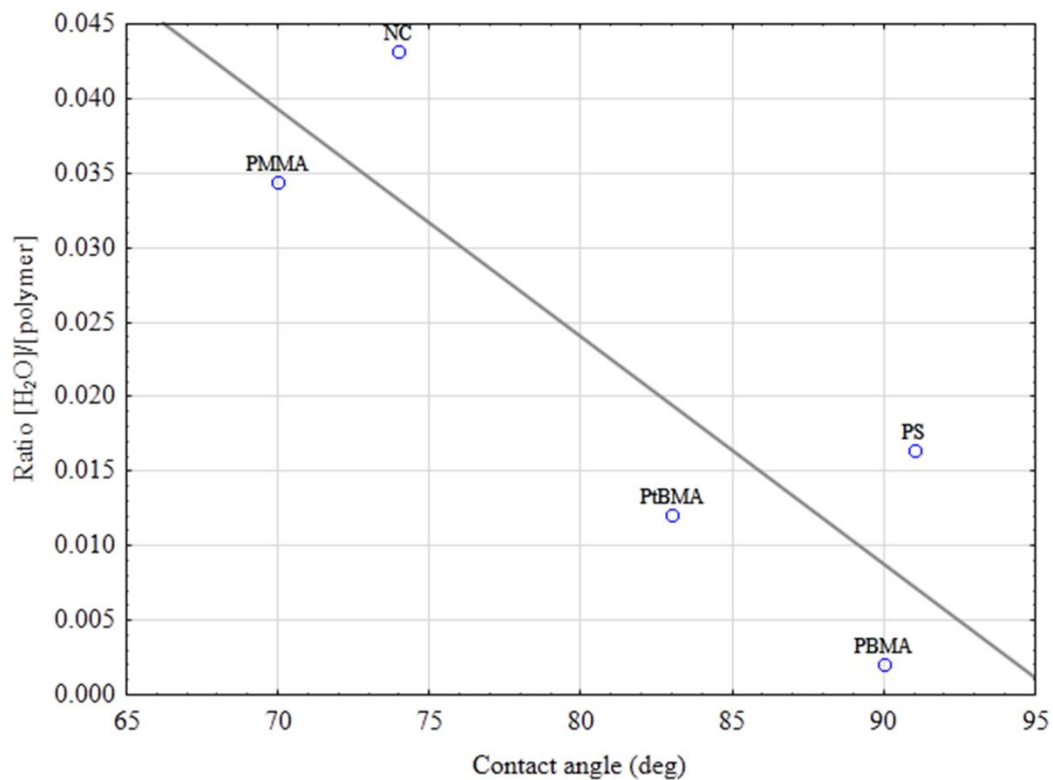


Figure SI1. Polymer water uptake vs. its hydrophobicity, expressed as contact angle.

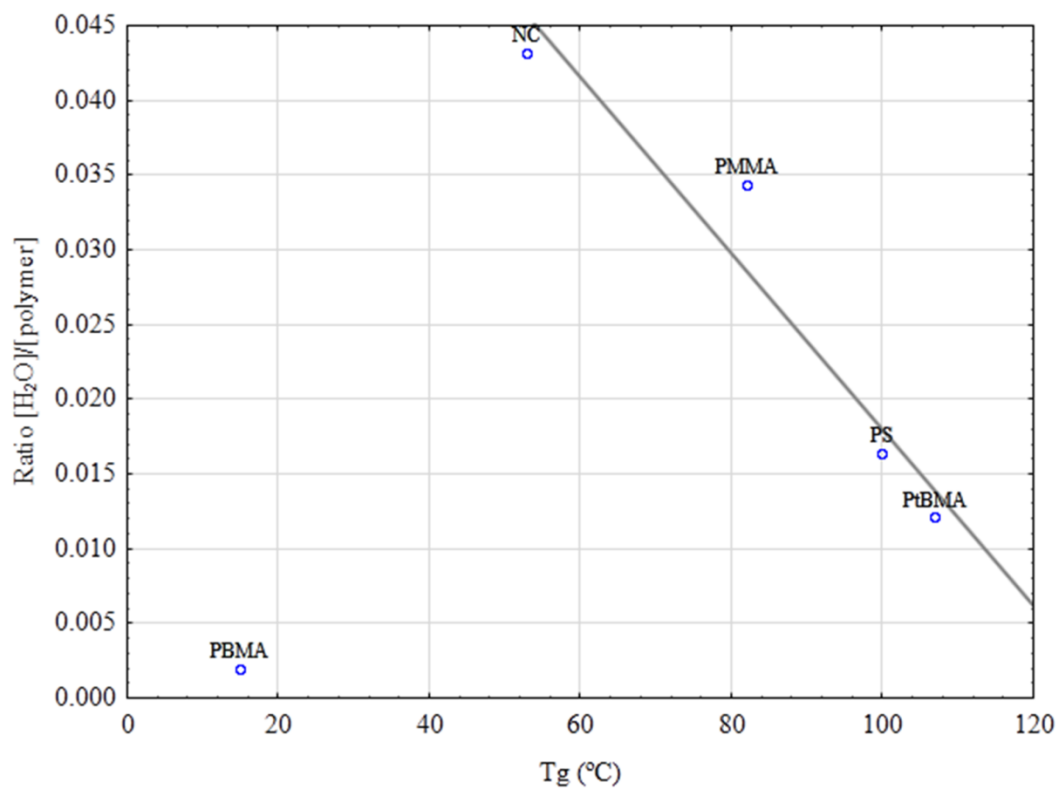


Figure SI2. Polymer water uptake vs. its glass transition (T_g) of the polymer.

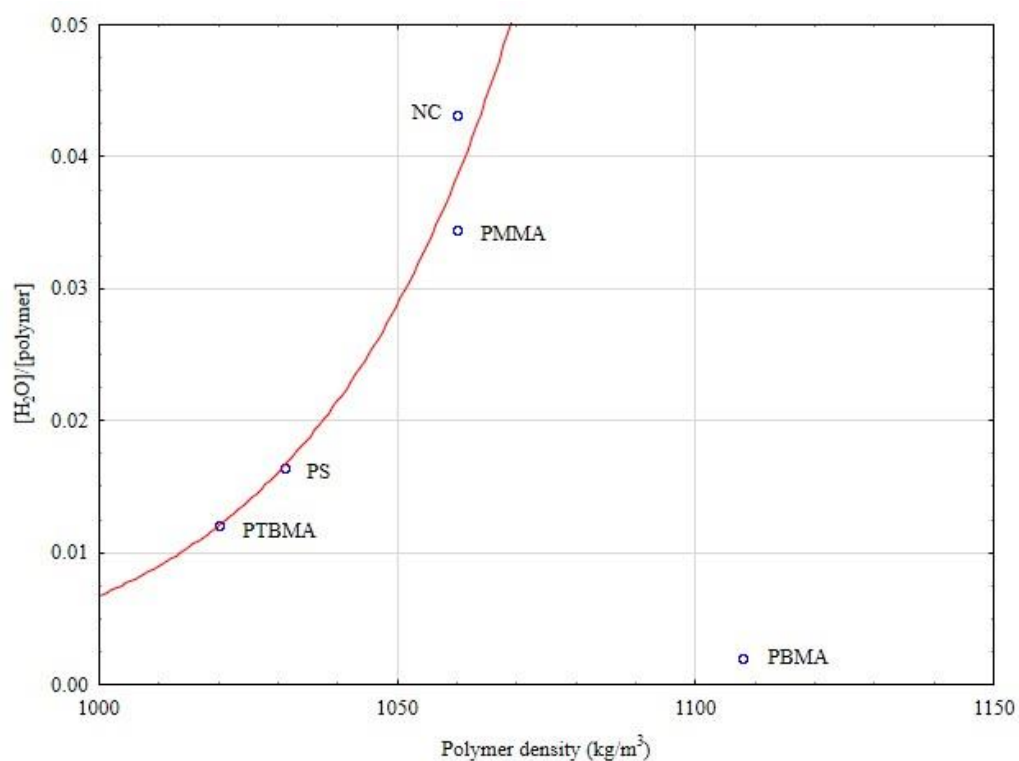


Figure SI3. Polymer water uptake vs. polymer density. Fit for glassy polymers only.

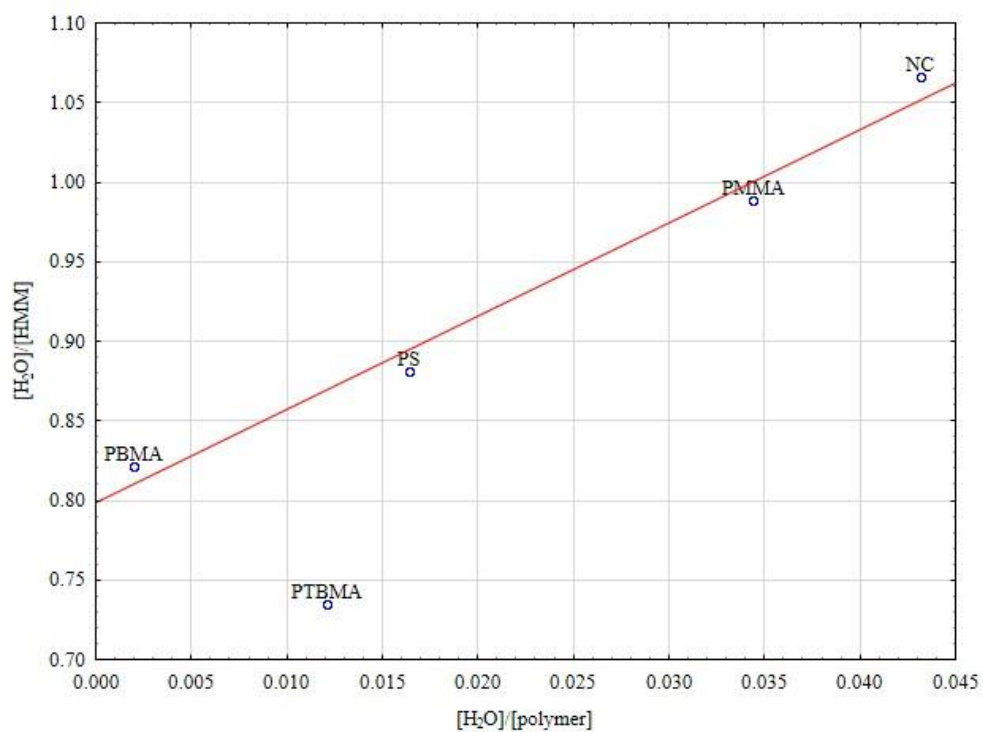


Figure SI4. HMM water uptake vs. polymer water uptake.

3. Proposed model of the polymer-HMM-liquid layers

It is reasonable to assume that the protein layer has a non-uniform density along its vertical cross-section. The HMM molecule is 80-90 nm long including the S1 heads (~19 nm long), with a diameter of 2-3 nm. Because the actual conformation of the adsorbed HMM molecule is not known, we considered that the molecule has a part of it bound to the surface and a part, containing one or both heads, oriented vertically. The bound part forms a layer several nm's thick, a protein rich layer with less adsorbed water; and one layer with a thickness around 15-35 nm, which is rich in water, and which is made of S1 heads of the HMM molecule, capable of interacting with the gliding filament. These layers are treated separately because they have different physical properties, i.e., densities and dynamical moduli. Because the first layer is very thin, it can be treated within the Sauerbrey approximation. The other layer, which is the one containing the heads, behaves like a gel, and therefore it is modeled as a viscoelastic layer. One restriction that needs to be imposed is that the number of molecules in each layer be the same (obviously this implies the assumption that, on average, each HMM molecule contributes to both layers.). Therefore, the HMM layer is divided into a more compact sub-layer with a thickness of several nm's; and a lower density layer, which is thicker, made of more vertically-orientated heads and containing water molecules. Each sub-layer is treated distinctly with regard to the contribution to the acoustic impedance. The overall contribution to the acoustic impedance is calculated by matrix multiplication, starting from the outmost layer, the HMM solution. The results of these calculations are presented in Table 2. One observation is that the density of the protein layer for all surfaces is approximately the same, i.e., $\sim 1090 \text{ Kg m}^{-3}$, which is close to the values reported elsewhere (e.g., 1040 Kg m^{-3} (Hook et al. 2002)). However, there are notable differences in the thickness of the protein layers, which suggests different conformations of the motor protein on different surfaces, and possibly its surface-bound site. In order to relate the number of molecules to the number of heads, we must estimate the proportion of the mass of the heads in the total mass of the gel layer. One possibility is to consider that, statistically, one head is part of the bottom layer and one head is part of the gel layer; and in this case we obtain the same number of heads as the number of molecules. The values obtained for the total mass adsorbed and the number of heads is in the correct order of magnitude as found in other experiments, e.g., 350 ng/cm^2 for myosin on glass.(Harada 1990)

This large amount of water trapped in the protein layer will make the calculation of the actual number of HMM per unit surface only indicative. In this context, a comprehensive study(Albet-Torres et al. 2007) used a high value of the correction factor (i.e., 10x) to calculate the surface density of HMM molecules, citing a previous comparative study which compared QCM and SAW measurements for mucin. However, in our work we derived the data from first principles, i.e., without using a correction factor.

References to Supplementary Information

- Albet-Torres, N., O'Mahony, J., Charlton, C., Balaz, M., Lisboa, P., Aastrup, T., Mansson, A., Nicholls, I.A., 2007. Mode of heavy meromyosin adsorption and motor function correlated with surface hydrophobicity and charge. *Langmuir* 23, 11147-11156.
- Banda, L., Alcoutlabi, M., McKenna, G.B., 2006. Errors Induced in Quartz Crystal Mass Uptake Measurements by Nongravimetric Effects: Considerations beyond the EerNisse Caution. *J Polym Sci Part B: Polym Phys* 44, 801-814.
- Benes, E., 1984. *J. Appl. Phys.* 56, 606.
- Dainty, M., Kleinzeller, A., Lawrence, A.S.C., Miall, M., Needham, J., Needham, D.M., Shen, S.C., 1944. Studies on the anomalous viscosity and flow-birefringence of protein solutions III. Changes in these properties of myosin solutions in relation to adenosinetriphosphate and muscular contraction. *Journal of General Physiology* 27(4), 355-399.
- Granstaff, V.E., Martin, S.J., 1994. Characterization of a thickness-shear mode quartz resonator with multiple nonpiezoelectric layers. *J. Appl. Phys.* 75(3), 1319-1329.
- Harada, Y., 1990. Mechanochemical coupling in actomyosin energy transduction studied by in vitro movement assay. *Journal of Molecular Biology* 216, 49-68.
- Hook, F., Kasemo, B., Nylander, T., Fant, C., Sott, K., Elwing, H., 2001. Variations in coupled water, viscoelastic properties, and film thickness of a Mefp-1 protein film during adsorption and cross-linking: A quartz crystal microbalance with dissipation monitoring, ellipsometry, and surface plasmon resonance study. *Analytical Chemistry* 73(24), 5796-5804.
- Hook, F., Voros, J., Rodahl, M., Kurrat, R., Boni, P., Ramsden, J.J., Textor, M., Spencer, N.D., Tengvall, P., Gold, J., Kasemo, B., 2002. A comparative study of protein adsorption on titanium oxide surfaces using in situ ellipsometry, optical waveguide lightmode spectroscopy, and quartz crystal microbalance/dissipation. *Colloids and Surfaces B-Biointerfaces* 24(2), 155-170.
- Johannsmann, D., 1999. Viscoelastic analysis of organic thin films on quartz resonators. *Macromol. Chem. Phys.* 200, 501-513.
- Kanazawa, K.K., Gordon, J.G., 1985. Frequency of a quartz microbalance in contact with liquid. *Analytical Chemistry* 57(8), 1770-1771.
- Lucklum, R., Behling, C., Hauptmann, P., 1999. Role of mass accumulation and viscoelastic film properties for the response of acoustic-wave based chemical sensors. *Analytical Chemistry* 71(13), 2488-2496.
- Lucklum, R., Behling, C., Hauptmann, P., 2000. Gravimetric and non-gravimetric chemical quartz crystal resonators. *Sensors and Actuators B-Chemical* 65(1-3), 277-283.
- Lucklum, R., Hauptmann, P., 2000a. The Delta f-Delta R QCM technique: an approach to an advanced sensor signal interpretation. *Electrochimica Acta* 45(22-23), 3907-3916.
- Lucklum, R., Hauptmann, P., 2000b. The quartz crystal microbalance: mass sensitivity, viscoelasticity and acoustic amplification. *Sensors and Actuators B-Chemical* 70(1-3), 30-36.
- Martin, S.J., Frye, G.C., Senturia, S.D., 1994. Dynamics and response of polymer-coated surface-acoustic-wave devices - Effect of viscoelastic properties and film resonance. *Analytical Chemistry* 66(14), 2201-2219.
- Mason, W.P., 1948. *Piezoelectric Crystals and their Applications to Ultrasonics*. Van Nostand, Princeton.
- Sauerbrey, G., 1959. Verwendung von schwingquartzen zur wagung dunner schichten und zur mikrowagung. *Zeitschrift Fur Physik* 155(2), 206-222.

Shen, D.Z., Huang, M.H., Chow, L.M., Yang, M.S., 2001. Kinetic profile of the adsorption and conformational change of lysozyme on self-assembled monolayers as revealed by quartz crystal resonator. *Sensors and Actuators B-Chemical* 77(3), 664-670.

Yang, M.S., Thompson, M., 1993. Multiple chemical information from the thickness shear mode acoustic-wave sensor in the liquid-phase. *Analytical Chemistry* 65(9), 1158-1168.